# ON THE FORMULATION OF FINITE ELEMENT

MODELS OF CAM MECHANISMS.

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#### SYNOPSIS

This paper describes a general procedure of kineto elastodynamic analysis of cam mechanisms based on the finite element approach. The present discrete technique can be utilized to provide various versions of finite element models of planar or spatial more complicated cam mechanisms. The procedure is introduced by utilizing the cam operated transfer mechanism found in Koster's Work<sup>(1)</sup>.

#### 1 - INTRODUCTION

The dynamic analysis of cam mechanisms as prefectly rigid systems has become increasingly inadequate, since the necessary prerequisities for setting up the vibration would not be satisfied. To improve the representation of dynamic behavior of a cam mechanism, various methods have been performed such as for example the methods by Matlhew and Tesar<sup>(3)</sup>, Eiss<sup>(4)</sup>, Bloom and Radcliffe<sup>(5)</sup>. There is however a common criticism to the previously mentioned analysis concerning the indiscrminate modelling technique. In reference (1), Koster gave a Kineto-elasto-dynamic analysis method for which the common drawbacks of the mentioned methods can be avoided, but still Koster's technique is sufficiently applicable for those simple cam mechanisms with low degrees of freedom and still many questions on the modelling of the cam shaft remain.

To permit a closer simulation of the dynamic behaviour of actual complicated cam mechanisms than was possible with simple models, the finite element approach has been utilized in the present analysis. Since this method (7,8) is an efficient tool for

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can operated transfer mechanism of figure (2) represented by 15 DF finite element model consists of a 12-DF simulated follower set and of a 3-DF simulated can set as shown.

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The follower set is modelled by connecting a series of links and each link may be simulated by one element performing a typical type of motion such as longitudnal, torsional, and or fleuxeral. With regard to the topology of the set, the individual elements are meeting at either pin or rigid joints. It is of interest to note that a series of dynamic models with lower degrees of freedom may be generated by eliminating the selected number of nodal generalized forces or of nodal generalized coordinates concerned the simulated follower set. Equivalently the sizes of mass and stiffness matrices are reduced either by the condensation of matrices or by the elimination of the particular number of rows and columns. These concepts will be visualized through the reduction of the simulated follower set to be as derived by Koster (1,2).

In the modelling of the flexible can set, the camshaft is represented by an assembly of flexural torsional beams interconnected at a rigid joint where as the cam element, (inertial element), is assumed to be lumped. Using the conditions of invariance of the kinetic and potential energies under coordinate transformation p, the mass matrix  $m_{e}$  and the stiffness matrix  $K_{e}$  of the e th element shown in figure 1, can be easly formulated.

The formulation of characteristic  $M_q$  and  $K_q$  matrices of the entire mechanism are then built up by adjoing the characteristic matrices of the cam and of the follower sets. The adjoing process is carried out by pre and post multiplication of the element characteristic matrices by the coupling matrix which represents the compatibility conditions through the nodal displacements. In that view the presence of coupling may be represented schematatically by a kinematic coupling set, (governed by the cam curve slope) as found in reference (1).

#### 3 - THE MODELLING OF THE CAM TRANSFER MECHANISM:

the analysis.

The simulated can mechanism is considered as a combination of the follower set formed from five structural elements and of the can set formed from three elements. Both sets are coupled by means of the coupler set which is simulated by a kinematic mechanism (1,2) as shown in figure (3).

In the uncoupled position of the follower set, the independent parameters  $q_{j}^{f}$   $(j = 1, 2, ..., l^2)$  are selected as the generalized coordinates. In the presence of coupling an auxiliarly dependent coordinate  $q^*$  may be also utilized for simulating the cam action, (the motion machined in the cam).

The generalized coordinates  $q_j^h$  (j = 1, ...,4) are utilized for describing the configuration of the rigidly supported cam set. Hereby the configuration of the simulated mechanism can be completely described by employing  $q_j$  (j = 1,2,...,16) generalized coordinates.

Figure (3) shows that the model of the entire mechanism consists of three submodels: follower, cam and coupling sets. In the modelling process the typical element is regarded as an elastic element, lumped masses and or rigid massless elements.

#### 3-1: Modelling of the Follower Set:

Follower set represented by that 12-DF model consists of the finite elements (1) up to (5) interconnected at two active pin joints (I, II) and two active rigid joints (III, IV). Taking into account the compatibility conditions, the location of the set of generalized coordinates  $q_1^f \dots q_{12}^f$ are selected as shown in figure (4). In the coupled position of the follower set, The auxiliarly generalized coordinate  $q^*$ , (which is utilized to simulate the cam action) is located at the proper distance (a) from the passive joint 0. A typical element (e) of mass  $\hat{f}$  is regarded an elastic element of longitudnal rigidity EA, torsional rigidity GJ and fleuxeral rigidicy EI, having a uniform cross-sectional area A across its length 1. The element mass and stiffness matrices can be shown (7,8) to be

$$\underline{\mathbf{m}}_{\mathbf{e}} = \frac{p}{\mathbf{G}} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix}, \quad \underline{\mathbf{K}}_{\mathbf{e}} = \frac{\mathbf{KA}}{1} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \text{ for } \mathbf{e} = 2 \dots (1)$$

$$\underline{\mathbf{m}}_{\mathbf{e}} = \frac{\mathbf{I}}{\mathbf{G}} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix}, \quad \underline{\mathbf{K}}_{\mathbf{e}} = \mathbf{GJ}/1 \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \text{ for } \mathbf{e} = 4 \dots (2)$$

$$\underline{\mathbf{m}}_{\mathbf{e}} = \frac{f^{2}}{420} \begin{bmatrix} 156 & 221 & 54 & -131 \\ 4p^{2} & 131 & -31^{2} \\ symmetric & 156 & -221 \\ & & 41^{2} \end{bmatrix}, \quad \underline{\mathbf{K}}_{\mathbf{e}} = \frac{\mathbf{EI}}{1^{3}} \begin{bmatrix} 12 & 61 & -12 & 61 \\ 41^{2} & -61 & 21^{2} \\ symmetric^{12} & -61 \\ 41^{2} \end{bmatrix} \text{ for } \mathbf{e} = 3, 5$$

$$\dots \dots (3)$$

In the coupled position of the follower set it may be convenient to subdivide the input link, (at the location of the auxiary coordinate  $q^*$ ), into two fleuxeral elements as shown in figure (5). The mass and stiffness matrices measured in a local system ( $\delta_1^2$ ,  $\delta_2^2$ ,  $\delta_3^2$ ,  $\delta_4^2$ ) can be synthesized by adjoing the characteristic matrices of two elements, using the elimination and condensation techniques, here as.

$$\underline{\mathbf{m}}_{1} = \mathcal{A} \begin{bmatrix} \frac{13}{35}\mathbf{l} + \frac{\mathbf{a}}{1680} & \frac{11}{210(b^{2}-a^{2}) - \frac{\mathbf{a}^{2}}{560}} & 9b/70 & -13b^{2}/420 \\ \mathbf{a}^{3}/560 + \frac{\mathbf{a}^{3}+b^{3}}{105} & 13b^{2}/420 & -b^{3}/140 \\ \\ \mathbf{symmetric} & 13b/35 & -11 & b^{2}/420 \\ & b^{3}/105 \end{bmatrix} \\ \underline{\mathbf{K}}_{1} = \begin{bmatrix} 12\mathbf{EI}/\mathbf{a}^{3} + 12\mathbf{EI}/b^{3} & 6\mathbf{EI}/b^{2} - 9\mathbf{EI}/\mathbf{a}^{2} & -12\mathbf{EI}/b^{3} & 6\mathbf{EI}/b^{2} \\ & 5\mathbf{EI}/\mathbf{a} + 4\mathbf{EI}/b & -6\mathbf{EI}/b^{2} & 2\mathbf{EI}/b \\ & 12\mathbf{I}/b^{3} & -6\mathbf{EI}/b^{2} \\ & 4\mathbf{EI}/b \end{bmatrix} \dots (4) \\ \mathbf{K}_{4} = \begin{bmatrix} \mathbf{symmetric} & 12\mathbf{I}/b^{3} & -6\mathbf{EI}/b^{2} \\ & 4\mathbf{EI}/b \end{bmatrix}$$

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In figure (5) the elements of the follower set are shown separated. Appropriate displacments are labelled on each, measured in the local and in the gobal coordinate systems. With  $\lambda = \cos \phi$  and  $\mathcal{M} = \sin \phi$ , a transformation matrix R of the typical element (e), figure (1), may be defined (6,7).

 $\mathbf{R}_{\mathbf{e}} = \begin{bmatrix} \lambda & \mathbf{M} & \mathbf{o} & \mathbf{o} \\ \mathbf{o} & \mathbf{o} & \lambda & \mathbf{H} \end{bmatrix} \quad \text{for } \mathbf{e} = 2,4 ,$ 

and

$$\mathbf{R}_{e} = \begin{bmatrix} -\mathcal{H} \ \lambda & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -\mathcal{H} & \lambda & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix} \quad \text{for } e = 1,3,5 \dots (5)$$

The following holds

S = R U

ž = R Ú

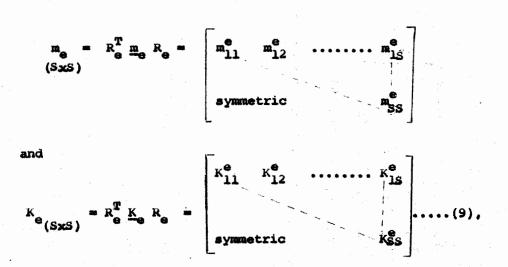
and

....(6)

Where  $\delta_{e}$  and  $U_{e}$  are sets of displacements at the two ends of the element in local and global coordinate systems as shown in figures (1) and (5). The invariance of kinetic and strain energies under coordinate transformation, such as:

and

facilitate the expressing the elemnt mass and stiffness matrices of the e  $\underline{th}$  element:



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Where S is the number of nodal displacements of the eth element. By the help of the previous equations, the symmetric element mass and stiffness matrices measured in a global system for the all isolated elements shown in figure (5) may be expressed as.

$$\begin{bmatrix} m_2 \end{bmatrix} = \frac{\gamma}{6} \begin{bmatrix} N_2 \end{bmatrix} , \begin{bmatrix} K_2 \end{bmatrix} = \frac{BA}{1^2} \begin{bmatrix} s_2 \end{bmatrix} ,$$

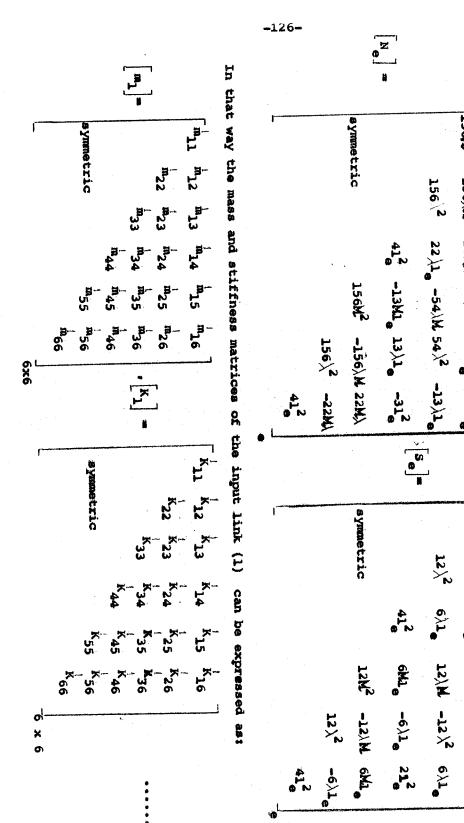
$$\begin{bmatrix} m_3 \end{bmatrix} = \frac{\gamma}{420} \begin{bmatrix} N_3 \end{bmatrix} , \begin{bmatrix} K_3 \end{bmatrix} = EI/1_3^3 \begin{bmatrix} s_3 \end{bmatrix} ,$$

$$\begin{bmatrix} m_4 \end{bmatrix} = \frac{I\rho}{6} \begin{bmatrix} N_4 \end{bmatrix} , \begin{bmatrix} K_4 \end{bmatrix} = \frac{GJ}{1} \begin{bmatrix} s_4 \end{bmatrix} ,$$

$$\begin{bmatrix} m_5 \end{bmatrix} = \frac{\gamma}{420} \begin{bmatrix} N_5 \end{bmatrix} , \begin{bmatrix} K_5 \end{bmatrix} = EI/1_5^3 \begin{bmatrix} s_5 \end{bmatrix} , \dots \dots (10)$$

Where the symmetric matrices [N] and [S] are defined:

$$\begin{bmatrix} N_{e} \end{bmatrix} = \begin{bmatrix} 2\lambda^{2} & 2\lambda^{M} & \lambda^{2} & \lambda M \\ 2N^{2} & M\lambda & M^{2} \end{bmatrix} \begin{bmatrix} S_{e} \end{bmatrix} = \begin{bmatrix} \lambda^{2} & \lambda M & -\lambda^{2} & -\lambda M \\ M^{2} & -\lambda M & -M^{2} \end{bmatrix}$$
 for e=2,4.  
symmetric  $2\lambda^{2} & 2\lambda M \\ 2M^{2} \end{bmatrix}$ 



... (11)

for e = 3, 5.

and

156M<sup>2</sup>

-156/H. -22M/ 54M<sup>2</sup> -54M1 13M1

12M2

-12/M -61 M -12M<sup>2</sup>

12)M -6MI

.

where the inertial coefficients  $m_{ij}$  and the elastic coefficients  $R_{ij}$  are given in the app endix.

Superimposing the kinetic and strain energies of the five elements shown in figure (5), the total energies may be expressed, using equations (7) and (8), as

and

 $v = \sum_{e=1}^{5} v_e = \frac{1}{2} v^T K_u v$  .....(13)

Here the mass  $M_u$  and stiffness  $K_u$  matrices of all elements comprising the follower set are derived by locating the element matrices of the five elements along the diagonal in the respective order (6). here has

and

and

$$K_u = \begin{bmatrix} K_1 \end{bmatrix} \dots \begin{bmatrix} K_5 \end{bmatrix}$$
  
24x24
(15)

The relationships between the generalized coordinates  $q^*$ ,  $q \frac{f}{1}$ , ...,  $q \frac{f}{12}$  and the nodal displacements  $U_1, U_2, \ldots, U_{24}$  of all elements of the follower set may be expressed as

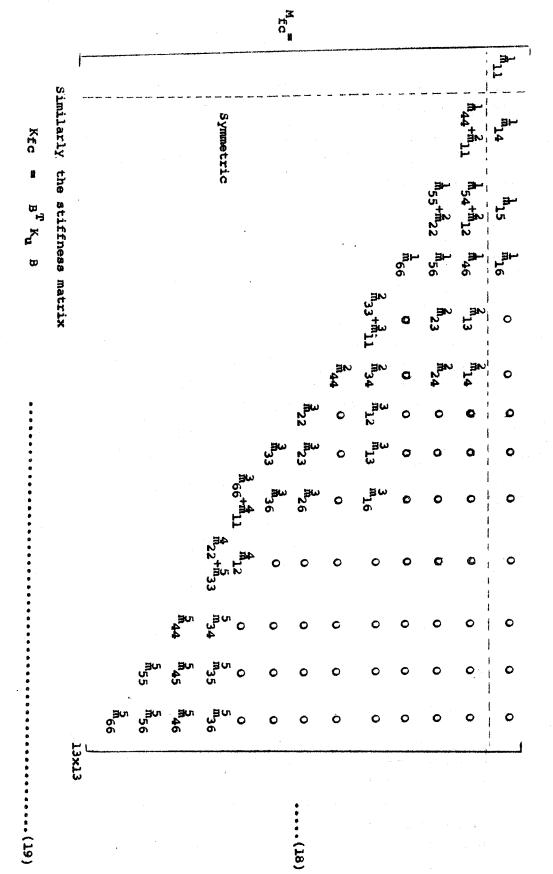
U = Bq .....(16)  $\dot{U} = B\dot{q}$  .....(17)

where B is the connecting matrix of order  $(24 \times 13)$ , which can be easly deduced according to the compatibility conditions through out the follower set.

Therefore the mass and stiffness matrices of the follower set in the coupled position may be expressed, using equations: (14) + (17):

 $M_{fc} = B^T M_U B$ 

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The expand form equation (19) is simillar to that given in equation (18) where the inertial coefficient  $m_{ij}^{e}$ and the elastic coefficient  $K_{ij}^{k}$  are given in the appendix.

In what follows, the derivation of relationship will be limited to the inertia properties, where as the elastic properties can be deduced similarly.

The (12 x 12) mass and stiffness matrices  $M_{fun}$  and  $K_{fun}$  of the follower set in the uncoupled position can be deduced by eliminating the first row and column, (corresponding to  $q^*$ ), from the matrices  $M_{fc}$  and  $K_{fc}$  as indicated by dashed lines in equation (18).

For many reasons it may be required to simulate the actual system with lower degrees of freedom model such as the shown 5-DF model of the follower set. In this case the (5 x 5) mass and stiffness matrices may be obtained by eliminating; the rows and columns which correspond to  $q^*, q_2^f$ ,  $q_3^f, q_5^f, q_6^f, q_7^f, q_{11}^f$  and  $q_{12}^f$ , figure (6), from the (13x13)  $M_{fc}$  and  $K_{fc}$  matrices.

The resulting matrices are:

N.B. The (5x5) stiffness matrix  $K_{f}$  has a similar form given in equation (20).

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The comparison between the matrices given in equation (20) and that derived in reference (1) shows that the developed matrices are more accurate for expressing the inertial and elastic properties, since the mutual effects of various types of deformations are still maintained.

A more precise representation of the inertial and elastical properties of the simulated follower set can be obtained by using the condensation technique (6) for the mass and stiffness matrices, since the kinematic compatibility conditions are still remain. For example the elimination of the nodal generalized forces  $Q^*$ ,  $Q_2$ ,  $Q_3$ ,  $Q_5$ ,  $Q_6$ ,  $Q_7$ ,  $Q_{11}$  and  $Q_{12}$  is equivalent to the reduction of the (13 x 13)  $M_{fc}$  and  $K_{fc}$ matrices as:

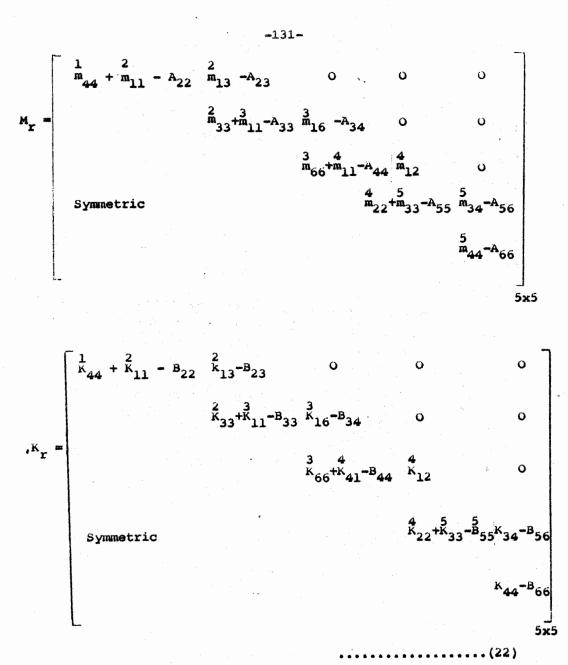
$$M_r = m_{11} - m_{12} = m_{22} = m_{21}$$
  
(5x5)

and

where

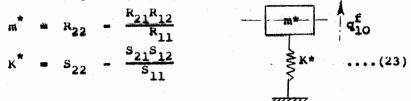
$$M_{fc} = \begin{bmatrix} m_{11} & m_{12} \\ (5x5) & (5x8) \\ m_{21} & m_{22} \\ (8x5) & (8x8) \end{bmatrix}, K_{fc} = \begin{bmatrix} K_{11} & K_{12} \\ (5x5) & (5x8) \\ K_{21} & K_{22} \\ (8x5) & (8x8) \end{bmatrix}$$

The expand forms of the reduced matrices may be expressed, using equations (18) (19) and (21):



Where  $A_{ij}$  and  $B_{ij}$  are functions of the mass and elastic coefficients of the individual elements in global coordinate system as shown in the appendix.

In that way, a more simplified dynamic model can be derived by further reducing the characteristic matrices. For example a single degree of freedom simulated follower set may be derived by recondensation of the  $M_r$  and  $K_r$  matrices. For example if  $q_{10}^f$ is selected to be the generalized coordinate, we have



It is seen that the form of equation (23) is more convenient for the dynamic investigation, since the parameters  $R_{11} \ldots S_{22}$  contain the inertial and elastical parameters of the elements comprising the follower set; (See the appendix).

## 3-2 Modelling of the Cam Set:

The rigidly supported cam set represented by that 4-DF model consists of three structural elements interconnected at three active nodes as shown in figure (7).

The camshaft is discretized into two fleuxeral-tersional elements 1 and 3 interconnected at the rigid joint II. The disk cam is idealized as nonelastic element lumped at the nodal point II as shown in figure (8).

Assuming uniform cross-section fleuxeral-torsional structural element. The characteristic matrices can be shown (6,7) to be

$$\begin{array}{c}
 \begin{bmatrix}
 13_{/35} & 0 & 0 & -13_{/35} & 0 & 0 \\
 13_{/35} & 0 & 0 & -13_{/35} & 0 \\
 I_{p/3A} & 0 & 0 & -I_{p/3A} \\
 Symmetric & 13_{/35} & 0 & 0 \\
 & & 13_{/35} & 0 \\
 & & & I_{p/3A}
\end{array}$$

$$\begin{array}{c}
 \begin{bmatrix}
 12 & 0 & 0 & -12 & 0 \\
 12 & 0 & 0 & -12 & 0 \\
 \psi 1^2 & 0 & 0 & -\psi 1^2 \\
 Sym. & 12 & 0 & 0 \\
 13_{/35} & 0 & 0 \\
 I_{p/3A}
\end{array}$$

Where  $\psi$  is a ratio GJ/EI.

Assuming the cam element of mass Y and of mass moment of inertia  $\mathcal{N}$  the mass matrix may be shown to be

With  $\lambda_{I} = \cos \phi_{I}$  and  $M_{I} = \sin \phi_{I}$  at the end I, the (6 x 6) transformation matrix may be defined as

$$\mathbf{R}_{e} = \begin{bmatrix} \mathbf{R}_{I} & \mathbf{0} \\ ----- & \mathbf{R}_{I} \\ \mathbf{0} & \mathbf{R}_{I} \end{bmatrix}, \quad \mathbf{R}_{I} = \begin{bmatrix} \lambda_{I} & \mathbf{M}_{I} & \mathbf{0} \\ -\mathbf{M}_{I} & \mathbf{2} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{1} \end{bmatrix}$$
 (26)

Therefore the  $(6 \times 6)$  characteristic matrices of the structural element may be expressed as:

2

	<sup>13</sup> /35 0 0	$\frac{13}{35}p = \frac{13}{35}p$	$\frac{13}{35}$ v 0		12 0 0	-12P	-12v	0]
	<sup>13</sup> /35 0	$-\frac{13}{35}v - \frac{13}{35}v$	<u>13</u> р о		12 0		-12P	0
	Ip/3A	0	0 -I <sub>p</sub> /3A		$\psi 1^2$	0	0 4	µ1 <sup>2</sup>
m <sub>e</sub> = -		13/35	1	Ke 13		 		
			0 0			12	0	0
	Symmetric	13	3/35 0		Symmetric		12	0
			Ip/3A		-		4	1 <sup>2</sup>

Where  $\mathbf{P} = \lambda_{\mathbf{I}} \lambda_{\mathbf{II}} + \mathcal{H}_{\mathbf{I}} \mathcal{H}_{\mathbf{II}} = \cos((\varphi_{\mathbf{II}} - \varphi_{\mathbf{I}}), \mathbf{v} = \lambda_{\mathbf{I}} \mathcal{H}_{\mathbf{II}} - \mathcal{H}_{\mathbf{I}} \lambda_{\mathbf{II}} = \sin(\varphi_{\mathbf{II}} - \varphi_{\mathbf{I}})$ 

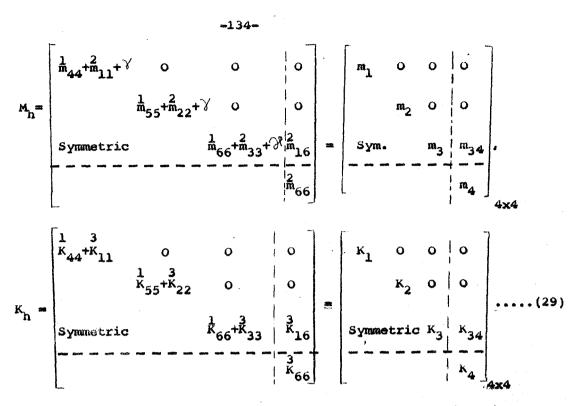
With the help of equations (12) to (15), the kinetic and strain engeries of the cam set are then given by

 $T = T_1 + T_2 + T_3$ ,  $V = V_1 + V_3$ ,

and the mass and stiffness matrices of isolated elements comprising the cam set are

$$M_{u} = \left[ m_{1} \right] \left[ m_{2} \right] \left[ m_{3} \right] \right], \quad K_{u} = \left[ K_{1} \right] \left[ 0 \right] \left[ K_{2} \right] \dots \dots \dots (28)$$

The relationships between the generalized coordinates  $q_1^h \oplus \oplus \oplus q_4^h$  and the nodal displacements  $U_1 \dots U_{18}$ , figure(8); may be expressed with the same forms given in equations (16) and (17). Therefore the mass and stiffness matrices of the cam set are then given by



Where the inertial and elastic coefficients are given in the appendix. If the inertia couple concerned with  $q_4^h$  is neglected, one get a statically and dynamically decoupling model similar to that derived in reference (1). The (3 x 3) mass and stiffness matrices are then obtained, as indicated by dashed lines in equation (29).

Refer to (1,2), the coupling between the follower and the cam sets is modeled kinematically a plane submechanism formed from a set of massless rigid links. The topology of the coupling set shown in figure (9) is governed by the transmission ratio(i) which is continuously variable and depends mainly on the cam curve slope (2).

The simple model derived in references (1, 2, 3, 4, 5) can be obtained as special cases of the simulated model developed here, such as for example the model derived by Bloom (5) can be deduced by neglecting the mass and flexability of the cam set in the vertical and tangential directions, whilst the neglection of the inertia and elastic parameters concerning the tangential and torsional direction leads to the Eiss's model (4) as shown in figure (10-a) and (10-b) respectively. It may be if interest to note that the inertial and elastical parameters  $(m_1, m_3, K_1, K_3)$  are expressed in terms of the inertial and elestical parameters of the elements comprising the cam set. Therefore the effect of local modifications on the dynamic characteristics of the system can be easly investigated.

3-3 Synthesis of the characteristic matrices of the simulated

### Cam-operated transfer Mechanism:

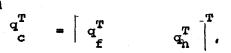
The characteristic matrices  $M_q$  and  $K_q$  of the entire mechanism are synthesized by the respective, pre and post multiplication of the characteristic matrices of the follower and cam sets in the uncoupled positions by a transformation matrix A. The latter matrix which specifies the compatibility condition between the two sets, relates the (16 x 1)  $q_{un}$  uncoupled vector with the (15 x 1)  $q_c$  coupled vector of the entire mechanism here as

 $q_{un} = A q_{c}$ 

where from definitions, we have

 $\mathbf{q}_{\mathbf{un}}^{\mathrm{T}} = \begin{bmatrix} \mathbf{q}^{*} & \mathbf{q}_{\mathrm{f}}^{\mathrm{T}} & \mathbf{q}_{\mathrm{h}}^{\mathrm{T}} \end{bmatrix}^{\mathrm{T}},$ 

and



The auxiliarly coordinate  $q^{n}$ , (which simulates the motion machined in the cam), and the  $q_{1}^{h}$ ,  $q_{2}^{h}$  and  $q_{3}^{h}$ , (concerned with the vertical, tangential and torsional deformations of the cam set as shown in figure 9), are related `(1) in the following matrix form

	a *	-	D	<b>G</b>			
Where	ã	is j	the	q <sub>h</sub> meshing	vector	given	by,
	D	-	(° 🛶	1 <u>i</u>	i ]		

The through inspection of eqns (30) + (32) reveals that the expanded forms of the transformation matrix may be given in the following partitioned scheme

$$A = \begin{bmatrix} 0 & D \\ -136- \\ I & 0 \\ -1 & -1 \\ I & 0 \\ -1 & -1 \\ 0 & I \end{bmatrix}$$

$$(33)$$

$$16 \times 15$$

With the help of equations (18) and (19) the both matrices may be partitioned in conformance with equation (32). Thus

$$M_{fc} = \begin{bmatrix} M_{11} & M_{12} & 0 \\ M_{12} & M_{fun} & 0 \\ 0 & 0 & M_{h} \end{bmatrix}; K_{fc} = \begin{bmatrix} K_{11} & K_{12} & 0 \\ - & - & - \\ K_{12}^{+} & K_{fun} & 0 \\ - & - & - \\ 0 & 0 & K_{h} \end{bmatrix} \dots (34)$$

Where  $M_{11}$  and  $K_{11}$  are the inertial and elastical parameters corresponding to  $q^2$ . Hereby the mass and stiffness matrices of the simulated entire mechanism may be expressed, using equatiosn (30) + (34), as

1

Similarly

$$K_{q} = A^{T} K_{fc}^{A} = \begin{bmatrix} K_{fun} & K_{12}^{T} D \\ - - - - - & - - - - - \\ D^{T} K_{12} & D^{T} K_{11} D + K_{h} \end{bmatrix} \dots \dots (36)$$

Where from definitions, we have

Γ

- 1)  $M_{fun}$  and  $K_{fun}$  are the (12 x 12) characteristic matrices of the follower set in the uncoupled position,
- 2)  $M_{h}$  and  $K_{h}$  are the (3x3) characteristic matrices of the cam set in the uncoupled position.
- 3)  $D^{T} M_{11} D = M_{11} D^{T} D$  is the diagonal submatrix of order (3%3), (since  $M_{11}$  is a scalar quantily).

4)  $D^{T}M_{12}$  is the (3 x 12) off diagonal submatrix, here as  $D^{T}M_{12} = \begin{bmatrix} B \\ 0 \\ 0 \end{bmatrix} 0 \begin{bmatrix} 0 \\ 0 \end{bmatrix} \cdots$ ...(37) where  $B = \begin{bmatrix} -\frac{1}{m_{12}} & -\frac{1}{m_{13}} & -\frac{1}{m_{14}} \\ \frac{1}{m_{12}} & \frac{1}{m_{13}} & \frac{1}{m_{14}} \\ \frac{1}{m_{12}} & \frac{1}{m_{13}} & \frac{1}{m_{14}} \end{bmatrix} \cdots$ ....(38)

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The developed partitioned scheme factitates markedly the formulation effort for synthesizing the characteristic matrix of the entire cam mechanism. Since the partitioned matrices given in equations 35 and 36 can be built up successively and therefore the required size of computer is considerably reduced.

For the sake of comparison between the present method and Koster's, the foundamental frequancy is computed (See appendix 2) for the various versions of simulated models of the cam operate transfer mechanism by using the bound formula (9). From the calculation, it is shown that the error, introduced by Koster technique relative to the present work lie within 5.56 to 6.83 percent.

#### CONCLUSION

The present approximate method attempts to provide a sufficiently simple and powerful tool for generating various versions in the modelling of the rigidly supported cam mechanism based on the finite element approach.

The proposed procedure avoids the drawbacks which arise in the applications of many classical methods such as methods mentioned here. The simple methods given in references (1, 2, 3, 4, 5)can be derived as special cases and as crude approximations of the developed method.

The formulation of the mass and stiffness matrices are introduced in such a way to render the problem tackable by limited capacity computer, which affect sharply economically in the computation effort for the machine design. This is because the characteristic matrices even in the condensed forms include the inertial and elastic parameters of all elements comprising the original system in the deterministic forms. However the computed result of the

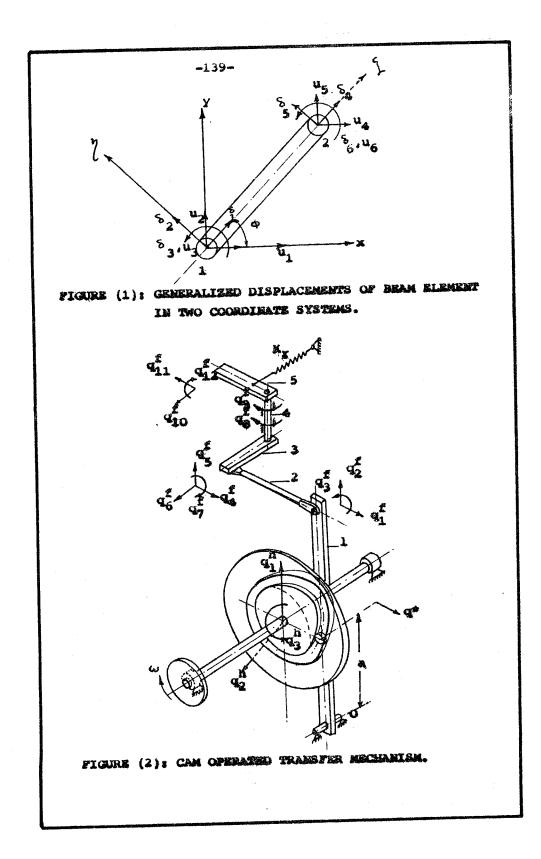
fundamental frequancy shows that present modelling technique is sufficient.

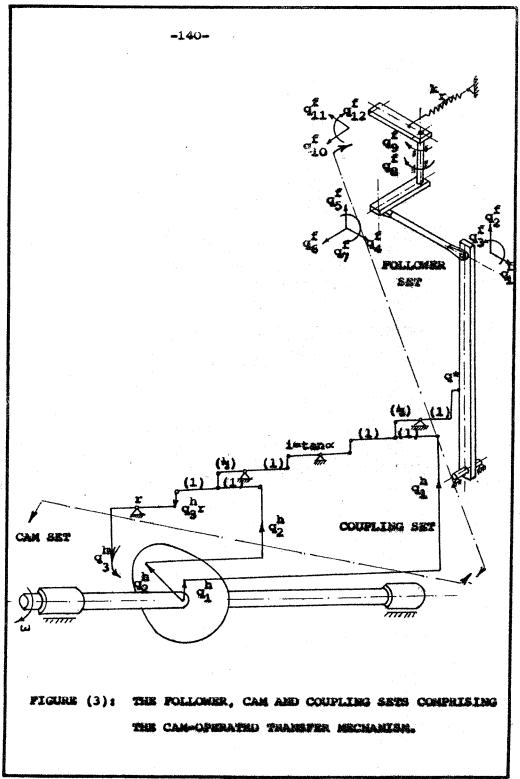
To integrate this work, further study and investigation concerning the variational effects of inertial and elastic parameters of actual constructive values of the system may be carried out in the future.

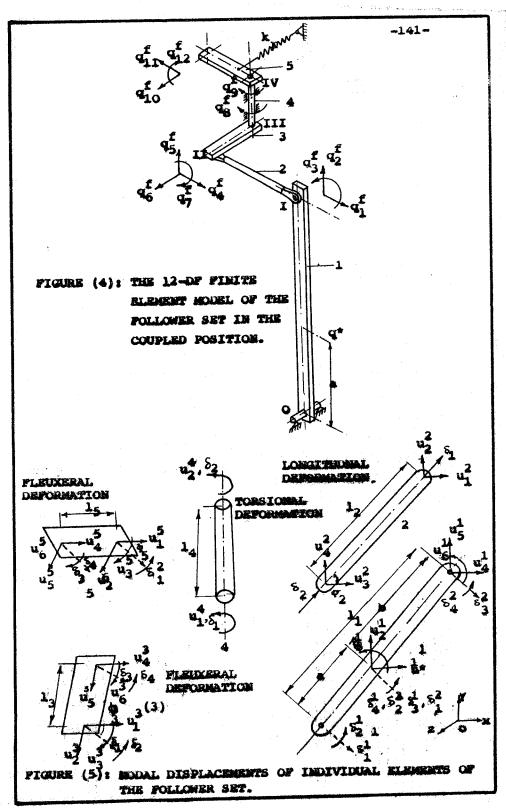
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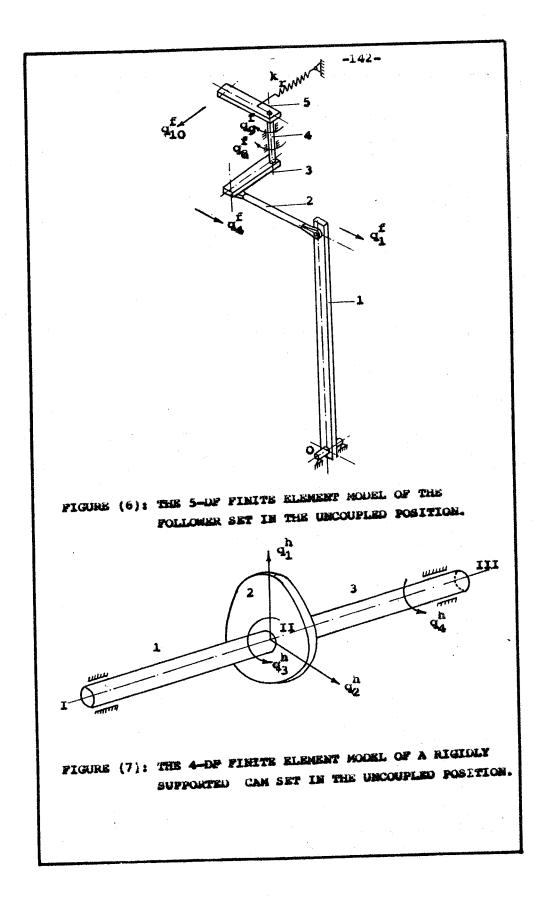
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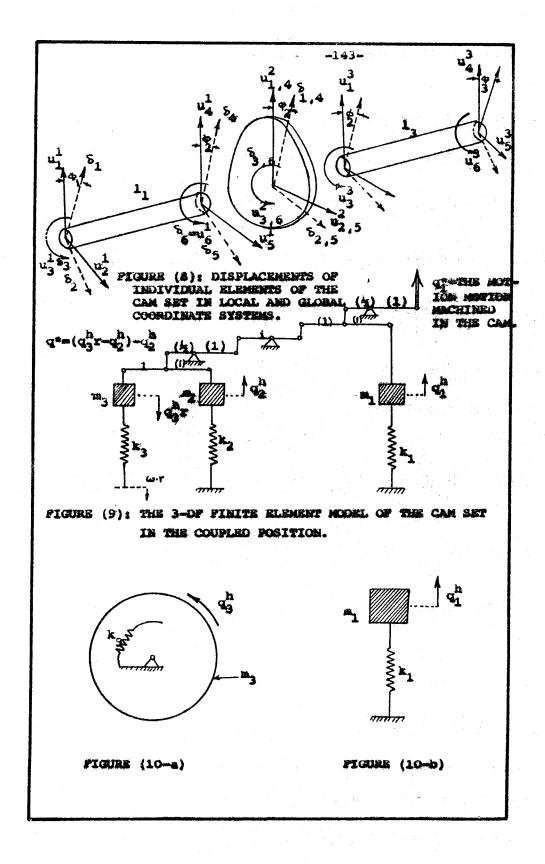
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## APPENDIX (1)

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The mass coefficients of the  $m_1$  matrix given in equation (11) are:

$$\frac{1}{m_{11}} = \int A(\frac{13}{35}1 + \frac{a}{1680}) \mathcal{H}^{2}; \frac{1}{m_{12}} = -\int A(\frac{13}{35}1 + \frac{a}{1680}); \frac{1}{m_{13}} = A(\frac{11}{210}(b^{2} - a^{2}) - \frac{a}{1680})$$

$$: \frac{1}{m_{14}} = \frac{9}{70} \int Ab\mathcal{H}^{2}; \frac{1}{m_{15}} = -\frac{9\mathcal{P}Ab}{70} \mathcal{H}\lambda; \frac{1}{m_{16}} = \frac{13\mathcal{P}Ab^{2}}{420} \mathcal{H}^{2}, \frac{1}{m_{22}} = \int A(\frac{13}{35}1 + \frac{a}{1680})\lambda^{2}$$

$$: \frac{1}{m_{23}} = \int A(\frac{11}{210}(b^{2} - a^{2}) - \frac{a^{2}}{560})^{2}; \frac{1}{m_{24}} = -\frac{9\mathcal{P}Ab}{70} \mathcal{H}\lambda; \frac{1}{m_{25}} = \frac{9\mathcal{P}Ab}{70} \mathcal{H}\lambda; \frac{1}{m_{26}} = \frac{13\mathcal{P}Ab^{3}}{420}\lambda$$

$$: \frac{1}{m_{33}} = \int A(\frac{a^{3}}{210} + \frac{a^{3} + b^{3}}{105}); \frac{1}{m_{34}} = -\frac{13\mathcal{P}Ab^{2}}{420} \mathcal{H}; \frac{1}{m_{35}} = \frac{13\mathcal{P}Ab^{2}}{420} \mathcal{H}, \frac{1}{m_{36}} = -\frac{\mathcal{P}Ab^{3}}{140};$$

$$: \frac{1}{m_{44}} = -\frac{13\mathcal{P}Ab}{35} \mathcal{H}^{2}, \frac{1}{m_{45}} = -\frac{13\mathcal{P}Ab}{35} \mathcal{H}\lambda; \frac{1}{m_{46}} = -\frac{11\mathcal{P}Ab^{2}}{420} \mathcal{H}; \frac{1}{m_{55}} = -\frac{13\mathcal{P}Ab}{35} \lambda^{2}, \frac{1}{m_{56}} = -\frac{11\mathcal{P}Ab}{420}\lambda$$

$$: and \frac{1}{m_{66}} = -\frac{\mathcal{P}Ab^{2}}{105}; \text{ where } a + b = 1$$

The elastic coefficients of the  $K_1$  matrix given in equation (11) are:

$$\begin{aligned} \mathbf{k}_{11} &= (\frac{12EI}{b^3} + \frac{21EI}{a^3}) \mathcal{H}^2, \mathbf{k}_{12} &= -(\frac{12EI}{b^3} + \frac{21EI}{a^3}) \mathcal{H}_3; \mathbf{k}_{13} &= (\frac{9EI}{a^2} - \frac{6EI}{b^2}) \mathcal{H}; \mathbf{k}_{14} &= -(\frac{12EI}{b^3}) \mathcal{H}^2 \\ &: \mathbf{k}_{15} &= (\frac{12EI}{b^3}) \mathcal{H}_3; \mathbf{k}_{16} &= -(\frac{6EI}{6^2}) \mathcal{H}^2; \mathbf{k}_{22} &= (\frac{12EI}{b^3} + \frac{21EI}{a^3}) \lambda^2; \mathbf{k}_{23} &= (\frac{6EI}{b^2} - \frac{9EI}{a^2}) \lambda \\ &: \mathbf{k}_{24} &= -(\frac{12EI}{b^3}) \mathcal{H}; \mathbf{k}_{25} &= -(\frac{12EI}{b^3}) \lambda^2; \mathbf{k}_{26} &= \frac{6EI}{b^2} \lambda; \mathbf{k}_{33} &= (\frac{4EI}{b} + \frac{5EI}{a}), \mathbf{k}_{34} &= \frac{6EI}{b^2} \mathcal{H} \\ &: \mathbf{k}_{35} &= -(\frac{6EI}{b^2}) \lambda; \mathbf{k}_{36} &= (\frac{2EI}{b}); \mathbf{k}_{44} &= (\frac{12EI}{b^3}) \mathcal{H}^2; \mathbf{k}_{45} &= -(\frac{12EI}{b^3}) \mathcal{H} \lambda \\ &: \mathbf{k}_{46} &= -(\frac{6EI}{b^2}) \mathcal{H}; \mathbf{k}_{55} &= \frac{12EI}{b^3} \lambda^2; \mathbf{k}_{56} &= -(\frac{6EI}{b^2}) \lambda; \mathbf{k}_{36} &= \frac{4EI}{b} \\ &: \text{The mass coefficients of the M}_{fc} matrix given in equation (18) are; \\ &: \mathbf{k}_{26} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda^2; \mathbf{k}_{26} &= \frac{2^2n}{b^2} &= 2^{2n} \mathcal{H} \lambda \\ &: \mathbf{k}_{35} &= -\frac{2^2n}{b^2} &= 2^2n \lambda \lambda^2; \mathbf{k}_{36} &= -\frac{2^2n}{b^3} &= 2^{2n} \mathcal{H} \lambda \\ &: \mathbf{k}_{46} &= -(\frac{6EI}{b^2}) \mathcal{H}; \mathbf{k}_{55} &= \frac{12EI}{b^3} \lambda^2; \mathbf{k}_{56} &= -(\frac{6EI}{b^2}) \lambda \\ &: \mathbf{k}_{66} &= \frac{4EI}{b} \\ &: \mathbf{k}_{66} &= \frac{4EI}{b} \\ &: \mathbf{k}_{66} &= \frac{4EI}{b} \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &= 2^2n \lambda \lambda \lambda \lambda \lambda \\ &: \mathbf{k}_{66} &= \frac{2^2n}{b^2} &$$

 $\hat{m}_{11} = \hat{m}_{33} = 2 \hat{m}_{13} = 2 \hat{p}_{A1} \lambda^2; \quad \hat{m}_{12} = \hat{m}_{34} = 2 \hat{m}_{14} = 2 \hat{m}_{23} = 2 \hat{p}_{A1} \mathcal{H} \lambda$   $\hat{m}_{22} = \hat{m}_{44} = 2 \hat{m}_{24} = 2 \hat{p}_{A1} \lambda^2; \quad \hat{m}_{11} = \frac{13}{35} \hat{p}_{A1} \mathcal{H}^2; \quad \hat{m}_{12} = -\frac{13}{35} \hat{p}_{A1} \mathcal{H} \lambda; \quad \hat{m}_{13} = -\frac{11 \hat{p}_{A1}^2}{210} \mathcal{H}$ 

Where;

$$\triangle = (\hat{\mathbf{h}}_{55} + \hat{\mathbf{h}}_{22}) F_{11} + \hat{\mathbf{h}}_{56} F_{12} + \hat{\mathbf{h}}_{24} F_{13}$$

$$F_{11} = \hat{\mathbf{h}}_{66} \cdot \hat{\mathbf{h}}_{44} \left[ \hat{\mathbf{h}}_{22} \, \hat{\mathbf{h}}_{33} - (\hat{\mathbf{h}}_{23})^2 \right] \left[ \hat{\mathbf{h}}_{66} \cdot \hat{\mathbf{h}}_{55} - (\hat{\mathbf{h}}_{56})^2 \right] : F_{12} = -\hat{\mathbf{h}}_{56} \frac{F_{11}}{\hat{\mathbf{h}}_{66}}$$

$$F_{13} = -\hat{\mathbf{h}}_{24} \frac{F_{11}}{\hat{\mathbf{m}}_{44}} : F_{22} = (F_{11})/\hat{\mathbf{h}}_{66} \cdot \hat{\mathbf{m}}_{44}) \left[ \hat{\mathbf{h}}_{44} (\hat{\mathbf{h}}_{55} + \hat{\mathbf{h}}_{22}) - (\hat{\mathbf{m}}_{24})^2 \right]$$

$$: F_{23} = (F_{11})/\hat{\mathbf{h}}_{66} \hat{\mathbf{m}}_{44}) (\hat{\mathbf{h}}_{56})^2 \cdot F_{33} = (F_{11})/\hat{\mathbf{h}}_{66} \hat{\mathbf{m}}_{44}) \left[ (\hat{\mathbf{h}}_{55} + \hat{\mathbf{h}}_{22}) \hat{\mathbf{h}}_{66} - (\hat{\mathbf{h}}_{56})^2 \right]$$

$$: F_{44} = \hat{\mathbf{m}}_{33} \left[ \left[ (\hat{\mathbf{h}}_{55} + \hat{\mathbf{h}}_{22}) \, \hat{\mathbf{h}}_{66} \hat{\mathbf{m}}_{44} - (\hat{\mathbf{m}}_{24})^2 (\hat{\mathbf{h}}_{66}) - (\hat{\mathbf{h}}_{56})^2 \, \hat{\mathbf{m}}_{44} \right] \left[ \hat{\mathbf{h}}_{66} \hat{\mathbf{h}}_{55} - (\hat{\mathbf{h}}_{56})^2 \right] \right]$$

$$: F_{45} = -\hat{\mathbf{m}}_{23} \frac{F_{44}}{\hat{\mathbf{h}}_{33}} \cdot F_{55} = \hat{\mathbf{m}}_{22} \frac{F_{44}}{\hat{\mathbf{h}}_{33}} : F_{66} = \hat{\mathbf{h}}_{66} \left[ \left[ (\hat{\mathbf{h}}_{55} + \hat{\mathbf{m}}_{22}) \hat{\mathbf{h}}_{66} \hat{\mathbf{m}}_{44} - (\hat{\mathbf{m}}_{24})^2 \right]$$

$$= \hat{\mathbf{h}}_{66} - (\hat{\mathbf{h}}_{56})^2 \, \hat{\mathbf{m}}_{44} \right] \left[ \hat{\mathbf{h}}_{33} \, \hat{\mathbf{h}}_{22} - (\hat{\mathbf{h}}_{33})^2 \right] \right]$$

$$: F_{67} = \hat{\mathbf{h}}_{56} \frac{F_{66}}{\hat{\mathbf{h}}_{66}} \text{ and } F_{77} = \hat{\mathbf{h}}_{55} \cdot \frac{F_{66}}{\hat{\mathbf{h}}_{66}}$$

The elastic coefficients of the 
$$K_r$$
 matrix given in equation (22) are:-  
 $B_{22} = \left[ \left[ k_{46} + k_{12} \right] L_{11} + k_{46} L_{21} + k_{14} L_{31} \right] \left( k_{45} + k_{12} \right] + k_{46} \left[ \left[ k_{45} + k_{12} \right] L_{12} + k_{46} L_{22} + k_{14} L_{32} \right] \right] + \left[ \left[ \left[ k_{45} + k_{12} \right] L_{13} + K_{46} L_{23} + k_{14} L_{33} \right] \left( k_{14} \right] \right] \cdot \frac{1}{\Delta} \right]$   
 $+ \left[ \left[ \left[ k_{45} + k_{12} \right] L_{13} + K_{46} L_{23} + k_{14} L_{33} \right] \left( k_{14} \right] \right] \cdot \frac{1}{\Delta} \right]$   
 $; B_{23} = \left[ \left[ \left[ k_{45} + k_{12} \right] L_{11} + k_{46} L_{21} + k_{14} L_{31} \right] k_{23} + \left[ k_{23} L_{13} + k_{34} L_{33} \right] \right] \cdot \frac{1}{\Delta} \right]$   
 $; B_{33} = \left[ \left[ \left( k_{23} L_{11} + k_{34} L_{31} \right) k_{23} + \left( k_{23} L_{12} + k_{34} L_{33} \right) k_{34} + \left( k_{12} L_{44} + k_{13} L_{45} \right) k_{12} \right]$   
 $+ \left( k_{12} L_{45} + k_{13} L_{55} \right) k_{13} \right] \cdot \frac{1}{\Delta}$ 

$$\begin{array}{l} \overset{-147-}{:}^{B_{34}} = \left[ (\hat{k}_{12}L_{44} + \hat{k}_{13}L_{45}) \, \hat{k}_{26} + (\hat{k}_{12}L_{45} + \hat{k}_{13} \, L_{55}) \, \hat{k}_{36} \, \right] \cdot \frac{1}{\Delta} \\ \overset{B}{:}^{B_{44}} = \left[ (\hat{k}_{26}L_{44} + \hat{k}_{36}L_{45}) \, \hat{k}_{26} + (\hat{k}_{26}L_{45} + \hat{k}_{36}L_{55}) \, \hat{k}_{36} \, \right] \cdot \frac{1}{\Delta} \\ \overset{B}{:}^{B_{55}} = \left[ (\hat{k}_{35}L_{66} + \hat{k}_{36}L_{67}) \, \hat{k}_{35} + (\hat{k}_{35}L_{67} + \hat{k}_{36}L_{77}) \, \hat{k}_{36} \, \right] \cdot \frac{1}{\Delta} \\ \overset{B}{:}^{B_{56}} = \left[ (\hat{k}_{54}L_{66} + \hat{k}_{46}L_{67}) \, \hat{k}_{35} + (\hat{k}_{45}L_{67} + \hat{k}_{46}L_{77}) \, \hat{k}_{46} \, \right] \cdot \frac{1}{\Delta} \\ \overset{A}{:}^{B_{56}} = \left[ (\hat{k}_{45}L_{66} + \hat{k}_{46}L_{67}) \, \hat{k}_{46} + (\hat{k}_{45}L_{67} + \hat{k}_{46}L_{77}) \, \hat{k}_{46} \, \right] \cdot \frac{1}{\Delta} \\ \overset{A}{:}^{B_{66}} = \left[ (\hat{k}_{45}L_{66} + \hat{k}_{46}L_{67}) \, \hat{k}_{46} + (\hat{k}_{45}L_{67} + \hat{k}_{46}L_{77}) \, \hat{k}_{46} \, \right] \cdot \frac{1}{\Delta} \\ \overset{A}{:}^{Hirst} \\ \overset{A}{:}^{Hirst} : L_{11} = \hat{k}_{66} \cdot \, \hat{k}_{44} \, \left[ \hat{k}_{22} \, \hat{k}_{33} - (\hat{k}_{23})^2 \right] \left[ \hat{k}_{66} \, \hat{k}_{55} - (\hat{k}_{56})^2 \, \right] : L_{12} = -\hat{k}_{56} \, \frac{L_{11}}{\hat{k}_{66}} \\ \vdots L_{13} = -\hat{k}_{24} \cdot \, \frac{L_{11}}{\hat{k}_{44}} \, \vdots L_{22} = (L_{11}/\hat{k}_{66} \cdot \, \hat{k}_{44}) \, \left[ \, \hat{k}_{44} ( \hat{k}_{55} + \, \hat{k}_{22}) - (\hat{k}_{24})^2 \, \right] \\ \vdots L_{23} = (L_{11}/\hat{k}_{66} \, \hat{k}_{44}) \, (\hat{k}_{56})^2 \, \vdots L_{33} = (L_{11}/\hat{k}_{66} \, \hat{k}_{44}) \, \left[ \, (\hat{k}_{55} + \, \hat{k}_{22}) \, \hat{k}_{66} - (\hat{k}_{56})^2 \, \right] \\ \vdots L_{45} = - \, \frac{L_{44} \cdot K_{23}^2}{\hat{k}_{33}} \, . \, L_{55} = \, \frac{L_{44}}{\hat{k}_{33}} \, \hat{k}_{22} \\ L_{66} = \hat{k}_{66} \, \left[ (\hat{k}_{55} + \, \hat{k}_{22}) \, \hat{k}_{66} \, \hat{k}_{44} - (\hat{k}_{24})^2 \, \hat{k}_{66} - (\hat{k}_{56})^2 \, \hat{k}_{44} \, \right] \left[ \hat{k}_{33} \, \hat{k}_{22} - (\hat{k}_{33})^2 \, \right] \\ \vdots L_{67} = \, \hat{k}_{56} \, \cdot \, \frac{L_{66}}{\hat{k}_{66}} \, \text{ and } L_{77} = \, \hat{k}_{55} \, \cdot \, \, \frac{L_{66}}{\hat{k}_{66}} \, . \\ \end{array}$$

 $R_{11} = (\frac{1}{m_{11}} - A_{11}) - \frac{1}{\Delta_2} \left[ (\frac{1}{m_{14}} - A_{11}) C_{11} - A_{13} C_{21} \right] (\frac{1}{m_{14}} - A_{21}) + A_{31} \left[ (A_{11} - \frac{1}{M_{14}}) C_{12} + A_{13} C_{22} \right]$ where  $1 \cdot 1 = \frac{1}{2} \cdot F_{-1}$ 1

$$A_{11} = \frac{1}{m_{15}} (\frac{1}{m_{15}} F_{11} + \frac{1}{m_{16}} F_{21}) + \frac{1}{m_{16}} (\frac{1}{m_{15}} F_{12} + \frac{1}{m_{16}} F_{22}) \cdot \frac{1}{\Delta}$$
  
$$A_{13} = (\frac{1}{m_{15}} F_{11} + \frac{1}{m_{16}} F_{21}) \frac{2}{m_{23}} + (\frac{1}{m_{15}} F_{13} + \frac{1}{m_{16}} F_{23}) \frac{2}{m_{34}} \cdot \frac{1}{\Delta}$$

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$$:^{A_{21}} = \left[ \left[ \left( \frac{1}{m_{45}} + \frac{2}{m_{12}} \right)^{F_{11}} + \frac{1}{m_{46}} + \frac{1}{m_{14}} + \frac{1}{m_{14}} \right]^{\frac{1}{m_{15}}} + \left[ \left( \frac{1}{m_{45}} + \frac{2}{m_{12}} \right)^{F_{12}} + \frac{1}{m_{46}} + \frac{2}{m_{14}} + \frac{2}{m_{14}} \right]^{\frac{1}{m_{16}}} \right]^{\frac{1}{m_{16}}} \\ :^{C_{11}} = \left( \frac{3}{m_{33}} + \frac{3}{m_{11}} - A_{33} \right) \left( \frac{3}{m_{66}} + \frac{4}{m_{11}} - A_{44} \right) \left( \frac{5}{m_{22}} + \frac{5}{m_{33}} - A_{55} \right) - \left( \frac{3}{m_{33}} + \frac{3}{m_{11}} - A_{33} \right) \left( \frac{4}{m_{14}} \right)^{2}$$

$$- (\overset{3}{m}_{16} - \overset{A}{}_{34})^{2} (\overset{5}{m}_{22} + \overset{5}{m}_{33} - \overset{A}{}_{55})$$

$$; c_{21} = (\overset{2}{A}_{23} - \overset{2}{m}_{13}) \left[ (\overset{3}{m}_{66} + \overset{4}{m}_{11} - \overset{A}{}_{44}) (\overset{5}{m}_{22} + \overset{5}{m}_{33} - \overset{A}{}_{55}) - (\overset{4}{m}_{12})^{2} \right] = c_{21} ; c_{22} = (\overset{1}{m}_{44} + \overset{2}{m}_{11} - \overset{A}{}_{11})$$

$$\left[ (\overset{5}{m}_{22} + \overset{5}{m}_{33} - \overset{A}{}_{55}) (\overset{3}{m}_{66} + \overset{4}{m}_{11} - \overset{A}{}_{44}) - (\overset{4}{m}_{12})^{2} \right] .$$

$$R_{12} = R_{21} = (\overset{A}{}_{56} - \overset{5}{m}_{34}) \left[ c_{14} (\overset{1}{m}_{14} - \overset{A}{}_{11}) - \overset{A}{}_{13} c_{24} \right] \overset{1}{\Delta}_{2}$$

Where;

$$c_{14} = \frac{4}{m_{12}} (A_{34} - \frac{3}{m_{16}}) (\frac{2}{m_{13}} - A_{23}) ; c_{24} = (\frac{1}{m_{44}} + \frac{2}{m_{11}} - A_{11}) (\frac{3}{m_{16}} - A_{34}) (\frac{4}{m_{12}})$$

$$R_{22} = (\frac{5}{m_{44}} - A_{66}) - \frac{1}{\Delta_2} (\frac{5}{m_{34}} - A_{56}) \left[ c_{44} (\frac{5}{m_{34}} - A_{56}) \right]$$
where;

$$c_{44} = (\frac{1}{m}_{44} + \frac{2}{m}_{11} - A_{22}) (\frac{3}{m}_{33} + \frac{3}{m}_{11} - A_{33}) (\frac{3}{m}_{66} + \frac{4}{m}_{11} - A_{44}) + (A_{22} - \frac{1}{m}_{44} + \frac{2}{m}_{11}) (\frac{3}{m}_{16} - A_{34})^{2}$$
$$+ (\frac{2}{m}_{13} - A_{23})^{2} (A_{44} - \frac{3}{m}_{66} - \frac{4}{m}_{11}); \bigtriangleup_{2} = (\frac{1}{m}_{44} + \frac{2}{m}_{11} - A_{22}) c_{11} + (A_{23} - \frac{2}{m}_{13}) c_{12}$$

The elastic coefficients of the  $k^*$  given in equation (23) are:-

$$s_{11} = (\hat{k}_{11} - B_{11}) - \frac{1}{\Delta_3} \left[ (\hat{k}_{14} - B_{11}) D_{11} - B_{13} D_{21} \right] (\hat{k}_{14} - B_{21}) + B_{31} \left[ (B_{11} - \hat{k}_{14}) D_{12} B_{13} D_{22} \right]$$
  
where;

$$B_{11} = \begin{bmatrix} k_{15} (k_{15}L_{11} + k_{16}L_{12}) + k_{16} (k_{15}L_{12} + k_{16}L_{22}) \end{bmatrix} \cdot \frac{1}{\Delta};$$

$$B_{13} = \begin{bmatrix} (k_{15}L_{11} + k_{16}L_{21}) k_{23} + (k_{15}L_{13} + k_{16}L_{23})k_{34} \end{bmatrix} \cdot \frac{1}{\Delta};$$

$$B_{21} = \begin{bmatrix} [(k_{46} + k_{12}) L_{11} + k_{46} L_{21} + k_{14} L_{31}] k_{15} + [(k_{45} + k_{12})L_{12} + k_{46}L_{22} + k_{46}L_{22} + k_{14} L_{31}] k_{15} + [(k_{45} + k_{12})L_{12} + k_{46}L_{22} + k_{46}L_{22} + k_{14} L_{31}] k_{15} + [(k_{45} + k_{12})L_{12} + k_{46}L_{22} + k_{46}L_{22} + k_{14} L_{31}] k_{15} + [(k_{45} + k_{12})L_{12} + k_{46}L_{22} + k_{46}L_{22} + k_{14} L_{31}] k_{15} + [(k_{45} + k_{12})L_{12} + k_{46}L_{22} + k$$

$$D_{21} = (B_{23} - \hat{k}_{13}) \left[ (\hat{k}_{66} + \hat{k}_{11} - B_{44}) (\hat{k}_{22} + \hat{k}_{33} - B_{55}) - (\hat{k}_{12})^2 \right] = D_{21}; \text{ and}$$

$$D_{22} = (\hat{k}_{44} + \hat{k}_{11} - B_{11}) \left[ (\hat{k}_{22} + \hat{k}_{33} - B_{55}) (\hat{k}_{66} + \hat{k}_{11} - B_{44}) - (\hat{k}_{12})^2 \right]$$

$$; s_{12} = s_{21} = (B_{56} - \hat{k}_{34}) \left[ D_{14} \hat{k}_{14} - B_{11} \right] - B_{13} D_{24} \right] \frac{1}{\Delta_3}$$

where;

$$D_{14} = \hat{R}_{12}(B_{34} - \hat{R}_{16})(\hat{R}_{13} - B_{23}); \text{ and } D_{24} = (\hat{R}_{44} + \hat{R}_{11} - B_{11})(\hat{R}_{16} - B_{34})(\hat{R}_{12})$$
  
and  
$$S_{22} = (\hat{R}_{44} - B_{66}) + \frac{1}{\lambda}(B_{56} - \hat{R}_{34}) \left[ D_{44}(\hat{R}_{34} - B_{56}) \right]$$

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$$\Delta_3$$
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where;  
 $D_{44} = (k_{44} + k_{11} - B_{22}) (k_{33} + k_{11} - B_{33}) (k_{66} + k_{11} - B_{44}) + (B_{22} - k_{44} - k_{11}) (k_{16} - B_{34})^2$   
 $+ (k_{13} - B_{23})^2 (B_{44} - k_{66} - k_{11}) \text{ and } \Delta_3 = (k_{44} + k_{11} - B_{22}) D_{11} + (B_{23} - k_{13}) D_{12}$ 

The mass coefficients of the  $M_h$  given in equation (29) are:

$$m_{1} = m_{2} = (\frac{13^{\rho}A1}{35}) + (\frac{13^{\rho}A1}{35})_{3} + m$$
  

$$m_{3} = (\frac{\rho_{1}p1}{3})_{1} + (\frac{\rho_{1}p1}{3})_{3} + 1 , \text{ and}$$
  

$$m_{34} = m_{43} = (\rho_{1}p1/3)_{3}$$

is the mass of cam disc. where m

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The elastic coefficients of  $k_{h}$  given in equation (29) are:

$$k_{1} = \left(\frac{12EI_{x}}{1^{3}}\right)_{1} + \left(\frac{12EI_{x}}{1^{3}}\right)_{3} ; k_{2} = \left(\frac{12EI_{y}}{1^{3}}\right)_{1} + \left(\frac{12EI_{y}}{1^{3}}\right)_{3}$$
  
$$k_{3} = \left(\frac{GI_{p}}{1}\right)_{1} + \left(\frac{GI_{p}}{1}\right)_{3} \text{ and } k_{4} = -k_{34} = \left(GI_{p}/1\right)_{3}$$

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## -150-

## Appendix (2)

The Calculation of the Fundamental Frequency

## I - Present Work :-

The kinetic energy of the follower set have 5 DF.  $(T) = \frac{1}{2} (1.4 57) \dot{q}_{1}^{2} + \frac{1}{2} \dot{q}_{1} \dot{q}_{4} + \frac{1}{2} (1.55714) \dot{q}_{4}^{2} + 0.069643 \dot{q}_{4} \dot{q}_{8}$   $+ \frac{1}{2} (0.53214) \dot{q}_{8}^{2} + 0.25 \dot{q}_{8} \dot{q}_{9} + \frac{1}{2} (0.53214) \dot{q}_{9}^{2} + 0.06964 \dot{q}_{9} \dot{q}_{10}$   $+ \frac{1}{2} (0.5572) \dot{q}_{10}^{2}$ 

The potential energy: -

$$(\mathbf{V}) = \frac{1}{2} (2.2819) q_1^2 - (2.0944) q_1 q_4 + \frac{1}{2} (5.64995) q_4^2 - 2.666 q_4 q_8$$
$$+ \frac{1}{2} (3.333) q_8^2 - 0.666 q_8 q_9 + \frac{1}{2} (3.333) q_9^2 - 2.666 q_9 q_{10}$$
$$+ \frac{1}{2} (3.555) q_{10}^2$$

Koster Work:-

$$(T) = \frac{1}{2} (1.6499) \dot{q}_{1}^{2} + \frac{1}{2} \dot{q}_{1} \dot{q}_{4} + \frac{1}{2} (1.3535) \dot{q}_{4}^{2} + (0.0589) \dot{q}_{4} \dot{q}_{8} + 0.52857 \dot{q}_{8}^{2} + 0.25 \dot{q}_{8} \dot{q}_{9} + \frac{1}{2} (0.5285) \dot{q}_{9}^{2} + (0.05286) \dot{q}_{9} \dot{q}_{10} + (0.5572) \dot{q}_{10}^{2} . (V) = \frac{1}{2} (2.20154) q_{1}^{2} - (2.0944) q_{1} q_{4} + \frac{1}{2} (2.9833) q_{4}^{2} - (.8888) q_{4} q_{8} + \frac{1}{2} (1.1852) q_{6}^{2} - (0.29629) q_{8} q_{9} + \frac{1}{2} (1.1852) q_{9}^{2} - 0.888 q_{9} q_{10} + \frac{1}{2} (0.888) q_{10}^{2}$$

The dynamic matrix of the follower set is given by:- $\begin{bmatrix} D \\ 5x5 \end{bmatrix} = \begin{bmatrix} K \end{bmatrix}^{-1} \begin{bmatrix} m \end{bmatrix}$  The largest eigenvalue is  $\lambda_1$  (fundamental frequency)  $\frac{T_r}{T_r} \frac{D^{k+1}}{D^k} \leqslant 1 \leqslant (T_r D^k)^{\frac{1}{k}}; T_r$  is trace of matrix[D]

The dynamic matrix [D] in present work is given by:-

	1.90933	1.62213	0.44116	0.325	0.1462	
	1.37099	1.528665	0.48065	0.35416	0.15895	
[D]=	1.21807	1.3814	0.4206	0.0169	0.1592	
- 3	0.6683	0.6263	0.0276	0.362	-0.338	
	0.4565	0.5178	0.0169	-0.170	1.00246	
	L				5	ix5

The dynamic matrix in Koster work is given by:-

					5	
	1.43647	1.40517	0.53852	0.52943	0.18606	
	1.61383	1.3306	J.55713	0.55625	0.19557	
	1.378548	1.2621999	0.24897	-0.161072	0.13092	
5x5	1.373668	1.17008	-0.1090875	0.386697	0.28679	
•	1.373918	1.258012	0.673965	1.2010614	2.06269	
					53	і к5

The dimension of the mechanism is taken as:-

11	=	a +	b	= 3 + 4	4 = 7				
1 <sub>2</sub>	-	3	cm.						
1 <sub>3</sub>	-	1.5	cm.			A =	I 277 =	constan	it
		1.5							
1 <sub>5</sub>	-	1.5	cm.						

In present Work

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 $T_r D^{16} = 1640228853$  $T_r D^{15} = 441501314.5$  Hence largest eigenvalue is  $\lambda_1$  (foundamental frequency)

$$\frac{1640228853}{441501314.5} \quad \lambda_1 \quad (441501314.5)^{\frac{1}{15}}$$

$$3.73253 \leqslant \lambda_1 \leqslant 3.7662$$

$$\therefore \lambda_1 = \frac{3.73253 + 3.7662}{2} = 3.749365 \quad 5^2$$

$$\therefore \text{ The foundamental frequency given by:}$$

$$\omega_1 = 0.5184414 \quad 5^{-1}$$
In Koster Work:-
$$T_r D^{16} = 6696135689$$

$$T_r D^{15} = 1601692192$$
The largest eigenvalue  $\lambda_1$  is given by:
$$\frac{6696135689}{1601692192} \leqslant \lambda_1 \leqslant (1601692192)^{\frac{1}{15}}$$

...4.21823 
$$\leqslant \lambda_1 \leqslant$$
 4.118077

$$\cdot \cdot \lambda_1 = \frac{4.21823 + 4.118077}{2} = 4.1681535 \text{ s}^2$$

s<sup>-1</sup>

.\*. The foundamental frequency is:

$$\omega_1 = 0.4896104$$

The relative deviation is given by:-

$$Z = \frac{0.5184414 - 0.4896104}{0.5184414} \times 100 = 5.56 \%$$

If the system represented as a single degree of freedom:-

1	In present		work:-	vork:-			
	<sup>K</sup> f	=	2.2819				
	<sup>m</sup> f	#	1.4857				
••	£	= /	$\frac{2.2819}{1.4857}$	. <b></b>	1.24173		

In Koster Work:-

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$$K_{f} = 2.20154$$

$$m_{f} = 1.6499$$

$$w_{f} = \sqrt{\frac{2.20154}{1.6499}} = 1.1551$$

$$\frac{7}{1.24173} = 1.1551 = 6.83\%$$

يهدف هذا البحث إلى تطبيق طريقة العناصر المحدود Finite element approach " للحصول على النبوزج الرياضي لآليات الكامات أخرى ما يمكن لالحقيقة والتحقيق هذا البحث اختيرت كامة تحرك تابسع ذو حركة تذيذبية ويحمل هذا التابع مجموعة من الأجسزاء لها جميع خواص الاهستزاز الطولى والمتلوى والمستعرض • ولتطبيق هذه الطريقة تم تقسيسم آليسة الكامة السي ثاسلات مجموعات وهي مجموعة التابع " Follower set" مشلة بمنظومة عديدة درجات الحرية ولها ١٣ درجسة حربة ومجموعة الكامة " Cam set " مثلة بمنظرمة ليا ثلاث درجات من الحربة والمجموعة التي تربط بين مجموعة التابع و مجموعة الكامة " Coupling set" وهي علاقة كينماتيكية تحتوى على جميع الخسواص الكينماتيكية لآليسة الكامة • أى أن في هذا البحث تم تمثيل آلية الكامسسسة يستظومة النها 17 درجة حربة والتخفيات هذا المسدد من درجات الحربة إلى أي عدد من المحاور الرام المرغوب فيهسا من درجات الحربة تم ذلك بر شريقتيسن أحدهما طريقة اخستزال Elemination technique" والطريقة الأخرى طريقة التكتيف "Condensation techinque" وقد وجدد أن طريق التكثيف أفضل من طريقة الاختزال • وقد وجد أن طريقة العناصر المحدودة مناسبسة لاتمثيسسل آليسات الكامات عن بعض الطرق المستخذمة من قبل وبالأخص في تمثيسل آليات المقعسدة وتسسم الدلك بمقاربتها ببعض الطرق التي استخدمها بعض الباحثيين السابقيين كما تمنساز أيضم بسبولة حساباتها باستخسدا، الحاسب العلى " Computor " الذي يؤدى إلى تخيسر الوقت والجهد في تصميم أجــزام الماكينات •